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Wilting Deformation of Leaves Using Cell Dynamics and Time-Varying External Forces

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Abstract

In this article, we proposed an algorithm to simulate 115 the withered deformation with wrinkle and curl of 116 plant leaves due to the dehydration, basing the cell 117 swelling force and time-varying external force. Firstly, 118 we proposed the leaf boundary expansion algorithm to 119 locate the feature points on the tips of the veins, and 120 calculated the discrete geodesic paths to construct the 121 primary veins. Secondly, we defined a novel 122 mass-spring system by the cell dynamics, and the 123 non-uniform mass distribution, to accelerate the 124 movement of the boundary cells. Thirdly, we defined 125 the cell swelling force on the osmosis, to generate the 126 wrinkle deformation by adjusting the dehydration and 127 the cell swelling force. Fourthly, we defined the system's 128 time-varying external forces on the feature points of the vein, to generate the curl deformation by adjusting the 129 initial values of the external forces and the multiple 130 iterative parameters. Finally, we used the implicit 131 midpoint method to solve the equation of motion 132 iteratively to achieve deformation simulation. The 133 experimental results show that our algorithm can 134 simulate the wrinkle and curl deformation caused by 135 dehydration and withering of leaves with higher 136 authenticity. 137

1. Introduction

The geometric modeling of plant leaves, and the simulation of leaf geometry deformation due to changes in ambient temperature and humidity, both are in strong demand in a variety of fields recently, e.g., in film and animation industry, video games, education, scientific research, especially in the agricultural informatization field [1, 2]. In this article, out fruitful work has resulted in a solution of simulating the wrinkle and curl deformation of leaves during the dehydration process.

Leaf veins are the mechanical support skeleton of leaves, which changes shape as the leaf curls and wrinkles. The correlative research work at home and abroad can be divided into interactive method [3], image deformation method [4], and biological drive method [5], etc. However, most of the previous studies of wilting deformation have not dealt with the mechanical supporting effect of veins, which leads to the of lack mechanical structural principles in their leaf models. As a result, the results of the simulations are not good enough and the computational efficiency is comparatively low.

One of the aims of this research has therefore been to improve the simulation authenticity of leaves curl and wrinkle deformation. In this article, we construct the primary veins of plant leaf by using the discrete geodesic path, and introduce a boundary expansion method to construct the tip feature points of the veins in a 3D leaf mesh model to achieve our goals. 174

Furthermore, Even though previous studies have made 175 great progress in this field, most of them ignored the cell 176 swelling force within the plant itself and the influence of 177 external factors, which are violate the principle of the cell 178 dynamics [6]. To solve these problems, first, we construct a 179 leaf mass-spring structure model according to the principle 180 of the cell dynamics; then we introduce a non-uniform mass 181 distribution to adjust the motion state of the cells as 182 plasmids in different regions to increase the deformation 183 acceleration of cells in the boundary region; moreover, we 184 introduce permeation effect of cells to define the cell 185 swelling force. The force can be controlled by adjusting the 186 dehydration rate to generate the wrinkle deformation; 187 finally, system's time-varying external forces are 188 introduced at the feature points of the vein tips to generate 189 the curl deformation by adjusting the initial value of the 190 time-varying external forces and the iterative parameters.

The main contributions and innovations of our work can 191 be summarized as follows: 192

• A leaf boundary expansion method was proposed to 193 locate the feature points on the tips of the veins. 194

• The primary veins are constructed by calculating the 195 discrete geodesic paths. 196

• A non-uniform distribution of leaf cell mass is defined 197 to optimize the deformation effect. 198

• Wrinkle deformation is effectively controlled by 199 introducing the cell expansion force.

• Curl deformation is adjusted by introducing the system's time-varying external forces.

2. Related work

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Changes in ambient temperature and humidity affect the water content of leaves, which in turn may lead to cell dehydration and cause the wilting deformation of the leaf wrinkles and curls. In this article, our simulation of the leaf wilting process can be divided into two parts: vein modeling and leaf cell mechanics modeling.

In terms of the veins modeling, in 2010, Qin and Xiao [3] published a paper in which they proposed an interactive modeling method that uses the user-interactive modification to create the 3D shape of leaf veins, which is highly dependent on the attributes provided by the users. However, this method has the disadvantage of ignoring the physiological characteristics of plants and the lack of biological and physical principles of the deformation process.

Recently, many studies have implemented more realistic simulations of leaf veins in 2D, e.g., in 2016, a contour-based leaf vein modeling method proposed by Kim and Kim [4] uses binary images of leaves to construct veins, and in 2017, biological driving methods (Runions et al. and Alsweis et al. [5, 7]) implemented the modeling and visualization of the leaf veins growing process. However, these simulations are only in the 2D plane and the latter two methods generated leaf edges are difficult to control.

225 In terms of the leaf deformation simulation, several 226 attempts have been made to improve the efficiency and 227 quality of deformation simulations. In 2009, A 228 physics-based two-layer leaf model proposed by Chi et al. [8] successfully simulates the deformation of leaves in 229 autumn, which upper layer represents the leaf sarcomeres 230 with higher shrinkage rates, and the lower layer represents 231 the vein cells that are more resistant to shrinkage and more 232 rigid. Two layers are connected to maintain the mechanical 233 stability of leaves. A mass-spring system is constructed 234 between those two layers to simulate the different kinetics 235 behaviors between the mesophyll cells and vein cells. As a 236 result, the two layers influence each other to jointly 237 determine the deformed shape of leaves. The advantage of 238 this method is that it better exploits biological features, but 239 the computational complexity of this method is high. 240 Similarly, in 2013, Jeong et al. [11] have proposed a 241 principle of the double-layered mass-spring model. The 242 difference is that in Jeong's method, both the vein and leaf 243 flesh models are double-layered, and a water penetration 244 model is constructed as well. As a result, this method can 245 adjust the degree of curl according to the differences in 246 water content between the various parts of the leaf, thus a 247 more realistic simulation effect can be obtained by the method. However, the veins and mesophyll are bi-layered, 248 which does not correspond to the objective physiological 249 structure of the leaf. Besides, the method has more masses and springs, which leads to the high computational complexity of the algorithm. Another method was proposed by Chi et al. [9] in 2009 to simulate the wilting and aging process of leaves based on the large-scale leaf sampling, however, the computational efficiency of this method is relatively low. 250 251 252 253

In order to simplify the deformation design process and 254 optimize the deformation simulation results, in 2015, Wei 255 et al. [12] proposed a half-process and half-physical leaf 256 deformation algorithm. The method first sets the target 257 state of the veins, and then constructs the intermediate state 258 sequence of the vein deformation by interpolating between 259 the initial state and the target state. The mesophyll is 260 deformed by the elastic force of the veins, which improve 261 the coordination between the spring coefficient and 262 deformation. However, the disadvantage is that the veins 263 model needs to be constructed manually, which is not 264 suitable for large-scale leaf modeling.

A novel method was proposed by Xiao and Chen [10] in 2011. This method takes the contraction strain field as the driving force of the wrinkling morphology, and controls the degree of leaf deformation by adjusting parameters. However, this method doesn't consider the role of veins, ignores structural analysis, and lacks biological principles. 270

Recently, in 2018, Liu et al. [13] proposed a method, 271 which is based on the feature that the wrinkle deformation 272 is more obvious and the curl deformation is weaker during 273 the initial stage of dehydration. When the water content 274 drops to a critical value, the shrinkage difference along the 275 boundary to the veins defines the curl deformation, which 276 provides some inspiration for the leaf deformation 277 algorithm in our article. 278

In view of all that has been mentioned so far, one may suppose that the construction of a realistic and efficient leaf vein model and the provision of an accurate and efficient deformation algorithm play a crucial role in the simulation of leaf wilt deformation.

283 In the section 3 of this article, we define the feature 284 points of the veins by using a leaf boundary expansion 285 algorithm, and calculate the discrete geodesic paths to 286 construct the primary veins; In section 4, we define the mass-spring system for cells, construct a non-uniform mass 287 distribution, give the calculation method of water content 288 under dehydration, define the cell swelling force and 289 time-varying external forces; Section 5 gives some 290 experimental results and analysis; Section 6 summarizes 291 the article. Figure 1 is the flow chart of the algorithm. 292

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3. Construct veins

In this section, in order to more accurately and efficiently generate the primary leaf veins, we propose a boundary expansion algorithm in the process of generating the primary leaf veins. The process is divided into three steps. We first construct a 3D mesh model of the leaves by the half-edge structure. Then, we use multi-dimensional scaling (MDS) to map the 3D leaf model into 2D space [14]. Once in 2D space, our boundary expansion algorithm is able to locate the local extreme points on the leaf boundary.

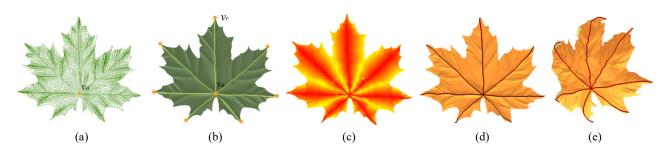


Fig. 1. Algorithm flow chart. In this article, we construct the 3D leaf model (a), locate the primary leaf veins (b), propose a non-uniform distribution of mass (c) to simulate the wrinkle deformation effect (d) and curl deformation effect (e) of plant leaf.

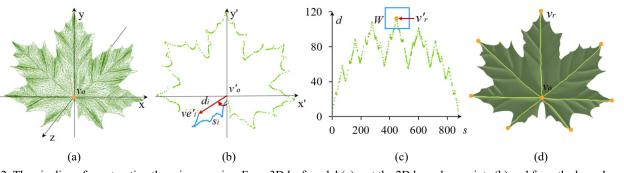


Fig. 2. The pipeline of constructing the primary veins. For a 3D leaf model (a), get the 2D boundary points (b) and form the boundary expansion coordinate (c) to construct the primary veins (d).

We shall take these points as the feature points of vein tips. Finally, we can use the obtained feature points in combination with discrete geodetic paths [15] to construct the primary leaf veins.

Let $G=\{V, E, F\}$ indicates the 3D mesh model of the leaf ,where $V = \{v_1, v_2, ..., v_n\}$ denotes the vertices set and $E = \{\langle e_{ii}, e_{ii} \rangle, v_i, v_i \in V\}$ is the set of symmetrical directed half-edges. When e_{ii} takes a non-zero value, it means that there exists a directed half-edge from vertex v_i pointing to vertex v_i . If half-edge e_{ij} has no opposite half-edge e_{ji} , i.e., e_{ji} is on the boundary. $V_E = \{ve_1, ve_2, ..., ve_{ne}\}$ denotes the set of such half-edges' endpoints, where ne is the number of boundary vertices. $F = \{f_1, f_2, ..., f_m\}$ is the set of the triangular faces of the leaf model.

 $D \in \mathbb{R}^{n \times n}$ is the distance matrix between the different339vertices of G, whose element gd_{ij} indicates the discrete340geodesic distance between the vertex v_i and v_j . Using MDS341to map the mesh model G from 3D space to 2D, the results342are denoted by G', the vertices set of G' is $\{v_1', v_2', ..., v_n'\}$.343Then the mapping process between the model G and G' can344be transferred to solve the following optimization [16]:

$$\min_{v_1,...,v_n} \sum_{i < j} \left(\left\| v'_i - v'_j \right\| - gd_{ij} \right)^2, i, j = 1, 2, \dots, n$$

For any vertex ve'_i at the boundary of the 2D mesh model G', d_i indicates the Euclidean distance between ve'_i and the petiole point v'_o (see Figur 2(b)). Calculating the cumulative length s_i of the boundary polyline from the boundary vertex ve'_i to the petiole point v'_o along the

boundary of G' in a clockwise direction. Taking the cumulative length s as the horizontal axis, d as the vertical axis and the (s_i, d_i) representing the coordinates of the boundary point ve'_i , the coordinate system of boundary expansion is drawn as shown in Figure 2(c) [17].

A sliding window *W* is defined with width *k* along the coordinate curve of the boundary vertices. In the window *W*, if there exists a boundary vertex v'_r with boundary vertices on both sides, satisfying the following conditions: $\exists r = i + k/2, d_i < ... < d_{r-1} < d_r < d_{r+1} < ... < d_{i+k}$, that boundary point v_r is defined as a feature point of the veins of model *G* in the 3D space.

In this article, we take the discrete measure path between 386 the petiole point v_0 and the feature point v_r of each vein on 387 the boundary as the primary veins of the leaf model *G*. The 388 primary veins and the feature points of veins shall affect the 389 non-uniform distribution of the initial mass of cells. And the feature points of veins are also the action points of 391 system's time-varying external forces 326

4. The wrinkle and curl deformation

In this section, in order to achieve an accurate simulation of wrinkle and curl deformation during leaf wilting process, we first define a plant leaf as a system of cells regarded as particles connected by springs according to cell dynamics [18] and mass-spring system [19]. Then we construct the mass-spring system according to the elastic force and damping force of the cells, generate the swelling force of leaves from the permeation effect of cells, and define the continuous time-varying external force at the feature points

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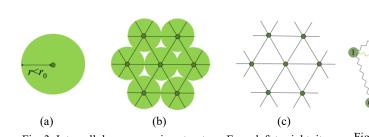


Fig. 3. Intercellular mass-spring structure. From left to right, it shows a single model of cell (a); a set of cell models (b); a simplified mass point-spring model (c).

 $v_{\rm r}$ of the veins. Finally, we combine all the above forces to implement the deformation simulation of leaves.

4.1. Mass-spring system

The cell dynamics-based model of plant leaf cells and its mass-spring system are shown in Figure 3. Figure 3(a) shows a single cell filled with cell sap, where r_0 and rrepresent the cell radius in the liquid-filled state of cells and the cell radius at shrinkage from dehydration, respectively. Figure 3(b) shows the mesophyll tissue formed by tightly arranged multiple cells, with plasmodesmata connected between the adjacent cells. Regarding cells as particles and plasmodesmata as springs, the mass-spring system corresponding to the mesophyll tissue of the multicellular leaf in Figure 3(b) is shown in Figure 3(c).

427 In this article, the vertices set V of the 3D mesh model 428 $G = \{V, E, F\}$ is used as the particles of the mass-spring system, and the edge set E represents the springs which 429 connecting the particles. The mass-spring system of G is 430 shown in Figure 4. Figure 4(a) shows the distribution model 431 of the spring. In the one-ring neighborhood of the particle v_i , 432 the number of particles connecting to v_i is m, and m 433 adjacent triangular patches are formed in the neighborhood. 434 In the mass-spring system, v_i is stretched or squeezed by m 435 particles in the one-ring neighborhood, the stress analysis 436 of v_i is shown in Figure 4(b). 437

4.2. Non-uniform distribution of mass

To better simulate the forces on each cell, in this article, we set up a non-uniform distribution of the initial mass of leaf cells corresponding to each particle in the mass-spring system. The mass m_i of a cell v_i is composed of the water content w_i and cytosol s_i . The initial mass of each mesophyll cell is obtained by the interpolation of the distance from the cell to the nearest leaf vein:

$$m_i^0 = m_{min} \left(1 - \left(1 - \frac{gd_i}{gd_{max}} \right)^{\mathsf{T}} \right) + m_{max} \left(1 - \frac{gd_i}{gd_{max}} \right)^{\mathsf{T}}$$

where $\tau \in [0,1]$ indicates the coefficient of mass distribution, when $\tau = 1$, the mass of cells decreases linearly

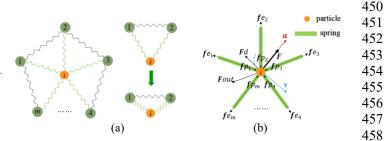


Fig. 4. Mass-spring model combined with cell mechanics. From left to right, it shows the adjacency deformation of particle v_i (a); the stress and motion state of particle v_i (b).

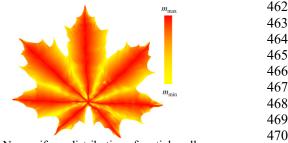


Fig. 5. Non-uniform distribution of particle cell mass

472 from vein cells to boundary cells, and when $\tau \neq 1$, the mass 473 of cells decreases nonlinearly. m_{\min} and m_{\max} are the 474 minimum and maximum values for the mass of the leaf vein cells. For each vein cell, there are multiple neighboring 475 mesophyll cells surrounding it. gd_{max} is the maximum value 476 of their distance from that vein cell. Figure 5 shows the 477 non-uniform distribution of the mass of the cells in a maple 478 leaf model, the initial mass of cells gradually decreases 479 from veins to boundaries. 480

4.3. Permeation and water content

483 The protoplast layer of plant cells can be regarded as a 484 permselective membrane, water molecules can move freely in and out of it, which is known as the permeation effect. 485 Due to the permeation between leaf cells, the concentration 486 of cytosol increases when dehydration occurs, after that, 487 water molecules enter the cell from outside to maintain the 488 balance of cytosol concentration. By using the Laplacian 489 operator of permeation which is improved by Jeong [11], 490 the variations of the water content during the withering and 491 dehydration process can be given as: 492

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$$w^{n+l} = (w^n - w_{loss}\Delta t)(1 - \Delta t L_i^n)^{-1}$$
 (1) 494
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where Δt indicates the iteration time step, w_{loss} denotes the 496 water content lost due to transpiration in unit time, i.e., the 497 dehydration rate. The larger the w_{loss} , the faster the cells 498 lose water. w_{loss} is defined by: 499

$$w_{loss} = THK$$
 (2)

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where T indicates the ambient temperature, H denotes the ambient humidity and K is the parameter for different leaves.

In Equation 1, L_i represents the Laplace operator of the current cell v_i , which can be obtained from the following equation:

$$L_{i}^{n} = \frac{2}{E} \sum_{j} \lambda \frac{b_{i} w_{j}^{n} - b_{j} w_{i}^{n}}{l_{ij} (m_{i}^{n} + m_{j}^{n})} , i=1,2,...,n, j=1,2,...,m$$

where *n* denotes the number of cells in the leaf model, i.e., the number of vertices in the 3D mesh model G; m is the number of cells in the one-ring neighborhood of v_i ; $E=\Sigma l_{ii}$, l_{ii} indicates the rest length of the spring; λ represents the diffusion coefficient of permeation; b_i represents the mass of the cytosolic of cell v_i ; w_i^n is the water content of the cell v_i , which can be obtained by: $w_i^n = m_i^n - b_i$.

4.4. Stress analysis

The mechanical properties of plant leaves are closely related to the mechanical behavior of single cells [20]. And the mechanical analysis of the cell is the key component to make the connection between the micro and macro biomechanics of the cells [21]. The macroscopic morphological changes of leaves are the manifestation of many factors in microscopic scales, such as the turgor pressure of each single cell, the penetration effect and the mechanical properties of the cell wall, etc. [22]

528 In order to obtain accurate and realistic force analysis 529 results, in this article, the movement of each cell v_i in the 530 leaf model G follows the laws of cell mechanics, which is 531 determined by its current mass m_i , spatial coordinate 532 position x_i , and the resultant force F_i , etc. As shown in 533 Figure 4(b), in addition to defining the elastic force Fe_i , and 534 the damping force Fd_i to maintain the stability of the 535 system, we also define the permeation expansion force Fp_i 536 of cells and the system's time-varying external forces Fo_i 537 acting on the feature points v_r of the veins. Therefore, the 538 resultant force F_i acting on the cells v_i can be obtained by: 539

$$F_i = Fe_i + Fd_i + Fp_i + Fo_i, i = 1, 2, ..., n$$

542 where $Fe_i = k_s \Sigma(|\mathbf{x}_{ij}| - l_{ij}) \mathbf{x}_{ij} / |\mathbf{x}_{ij}|, i=1, 2, ..., n, j=1, 2, ..., m,$ 543 which indicates the resultant force of the elastic at v_i from m 544 cells in the one-ring neighborhood of v_i , where k_s denotes 545 the coefficient of elastic; l_{ii} is the initial length of the spring 546 between the cell v_i and v_i ; $|\mathbf{x}_{ij}|$ is the length of the spring at 547 the current moment. The introduction of system damping force Fd_i avoids the problem of system instability caused 548 by excessive oscillation from the movement of the particles. 549 Fd_i is given by: $Fd_i = k_d \Sigma(v_i - v_i), i=1, 2, ..., n, j=1, 2, ..., m$ where k_d denotes the damping coefficient, v_i and v_j are the current velocity of the corresponding particle (see Equation 3).

550 Furthermore, the liquid inside the cell is wrapped by the 551 cell wall and the plasma membrane, generating pressure. 552 The difference between this pressure and the external 553 standard atmospheric pressure is defined as the swelling 554 force. The cell swelling force Fp_i acting on the particle v_i is 555 defined as: $Fp_i = k_p \sum \alpha_{ij} \Delta p_i x_{ii} / |x_{ij}|, i=1, 2, ..., n, j=1, 2, ..., m,$ 556 where k_p indicates the coefficient of swelling force; α_{ij} 557 denotes the area of the *j*-th triangle among the *m* triangles 558 formed by the particle cell v_i with the *m* cells in the one-ring 559 neighborhood; ΔP_i is the difference between the internal 560 and external pressure of the particle v_i at the current 561 moment; The direction of the swelling force is as same as 562 the direction where the particle v_i in the neighborhood of 563 the one-ring points to the current particle v_i . Considering 564 that during the leaves withering process, the occurrence of 565 permeation and dehydration shall lead to a decrease in the volume of cells. Set the initial intracellular pressure p_i^0 566 equals to the extracellular pressure p_{out} , both are standard 567 atmospheric pressure, then p_i can be updated iteratively 568 by: $p_i^{n+1} = p_i^n - (E_v w_{loss}) / (w_i^n \Delta t), i = 1, 2, ..., n$, then ΔP_i can be 569 570 updated by: $\Delta p_i^{n+1} = p_{out} - p_i^{n+1}$, and the elastic modulus E_v 571 can be obtained by $E_v = c / (1-2\mu)$. Where c denotes the cell 572 constant; μ indicates the Poisson 's ratio; w_i^n represents the 573 water content of the cell v_i at the current moment (in 574 Equation 1) and w_{loss} is the dehydration rate (in Equation 2).

575 In this article, the system's time-varying external force defined at the feature point v_r of the vein is given as: 576 $fout_r^{n+1} = fout_r^n + k_r \Delta t$, where k_r is a vector of adjustment 577 coefficient. Then the resultant force of the system's 578 time-varying external forces acting on cell v_i at each feature 579 point v_r can be obtained by: $F_{Oi} = \Sigma fout_r^{n+1} e^{-K_r g d_{ir}}, i=1, 2, ...,$ 580 n, r=1, 2, ..., q, where q is the number of the feature points 581 of the veins; gd_{ir} denotes the discrete geodesic distance 582 from the current particle v_i to the feature point of vein v_r ; K_r 583 indicates the corresponding adjustment coefficient, which 584 controls the influence degree of each feature point on the 585 cell v_i . 586

4.5. Numerical solution

Considering that the forward Euler method has low 589 complexity, but with a restricted time step. The iterative 590 process of this method shall collapse if the time step 591 exceeds the critical value, which makes the method less 592 stable. On the contrary, the backward Euler method has greater stability, but the computational complexity is higher.

Combining the above analysis, in this article, we choose a compromise between the above two methods by solving 597 the motion equation of the leaves' dehydration deformation 598 iteratively via the implicit midpoint method [23]. The 599 improved formulas for the velocity and displacement of the cells v_i are as follows:

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	Maple leaves	Poplar leaves	Tilia leaves	Ivy leaves	Willow leaves	Grape leaves
$Mass[m_{\min}, m_{\max}]$	[0.70, 1.50]	[0.75, 1.55]	[0.85, 1.65]	[0.8, 1.0]	[0.75, 1.55]	[0.80, 1.50]
Diffusion coefficient λ	0.1	0.1	0.1	0.1	0.1	0.1
Temperature T	20°C	20°C	20°C	20°C	20°C	20°C
Humidity <i>H</i>	50%	50%	50%	50%	50%	50%
Leaf parameter K	2.0×10^{-3}	4.0×10^{-3}	2.0×10^{-3}	4.0×10^{-3}	5.0×10^{-3}	4.0×10^{-3}
Dehydration rate w_{loss}	0.010	0.020	0.010	0.020	0.025	0.020
Spring coefficient k_s	4.0	4.0	4.0	4.0	4.0	4.0
Expansion coefficient kp	1.0	1.0	1.0	1.0	1.0	1.0
Cell constant <i>c</i>	2.0×10^{-4}					
Poisson's ratio μ	0.499	0.499	0.499	0.499	0.499	0.499
Damping coefficient k_d	0.1	0.1	0.1	0.1	0.1	0.1
External force coefficient K_r	1.0	1.0	0.5	1.0	0.8	1.0
Time interval Δt	0.1	0.1	0.1	0.1	0.1	0.1

Table 1. Parameters of the withering deformation

$$\mathbf{v}_{i}^{n+1} = \mathbf{v}_{i}^{n} + \Delta t F\left(\frac{\mathbf{x}_{i}^{n+1} + \mathbf{x}_{i}^{n}}{2}\right) / m_{i}^{n}, i = 1, 2, \dots, n$$
(3)

$$\mathbf{x}_{i}^{n+1} = \mathbf{x}_{i}^{n} + \frac{\Delta t}{2} \left(\mathbf{v}_{i}^{n+1} + \mathbf{v}_{i}^{n} \right), i = 1, 2, \dots, n$$
(4)

where m_i^n denotes the mass of v_i at the current moment. Due to the dehydration from cells permeation, the water content of the cells is decreasing with time. As a result, the mass of the cells decreases as the number of iterations increases.

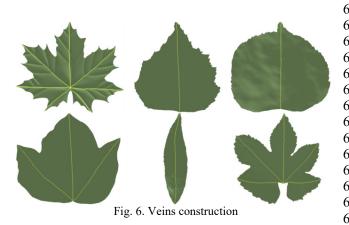
5. Experimental results and analysis

In this article, the experimental environment is a personal computer with Intel(R) Core(TM) i3-4160 CPU, Intel(R) HD Graphics 4400 GPU, 32GB of memory, 64-bit Windows operating system, and with Microsoft Visual Studio 2010 C++ and OpenGL as the environment of programming language. Under this condition, we have constructed the primary veins and simulated the wilting deformation of 6 species of plant leaves, including maple, poplar, tilia, willow, ivy and grape.

634 Overall, the leaf wilting deformation simulation process 635 can be decomposed into two parts: wrinkle deformation and 636 curl deformation. Upon the basis of the mass-spring system, 637 the result of wrinkle deformation is determined by cell 638 swelling force generated by the permeation of cellular 639 dehydration, and the curl deformation result depends on the 640 time-varying external force defined at the feature points of each vein. The leaves wilting deformation parameters of the experiment are provided in Table 1. 642

5.1. Veins construction

645 In this article, we first locate the feature points on the tips 646 of the veins by using the leaf boundary expansion algorithm, 647 and then construct the primary veins of the 3D leaf mesh 648 model with the petiole points based on the discrete geodesic 649 paths. The results of the primary veins we constructed are shown in Figure 6, there are 6 kinds of plant leaves, including poplar, linden, ivy, willow, and grape.



Obviously, in all the above species of leaves, the feature points of leaf veins can be accurately found, And the location and shape of the primary leaf veins closely matches the ones in the pictures of the real leaves. These results suggest that the boundary expansion algorithm we proposed in this article is well suited to high quality construction of veins

5.2. The effect of deformation

Figure 7 compares the deformation results of the leaves 689 under different forces. The first row of Figure 7 shows the 690 effect of deformation in different leaves without cell 691 swelling force and time-varying external force. The second row shows the effect when leaves are only affected by cell 692 swelling force and without time-varying external force, to 693 compare with the first row, the cell swelling force makes 694 the wrinkling effect of leaves much more pronounced. The 695 third row shows the result of deformation under the 696 influence of both cell swelling force and time-varying 697 external forces. Obviously, compared with the second row, 698 the curling of leaves is much more pronounced under the 699 time-varying external forces.

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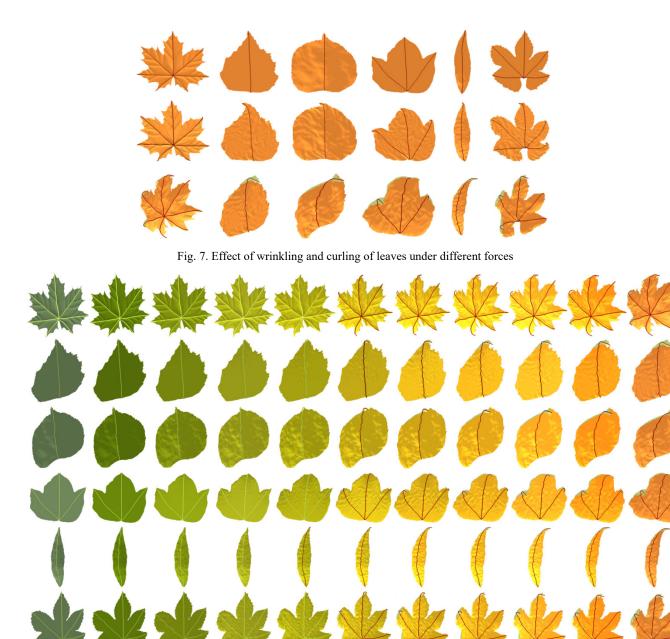


Fig. 8. The middle processes of the wrinkle and curl deformation

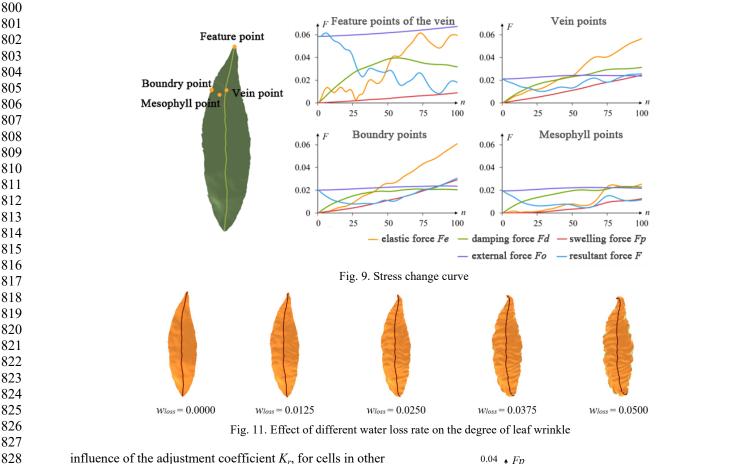
The above comparison shows that the cell swelling force we proposed in this article can be effectively used to generate the leaf wrinkle deformation, and the time-varying external forces are effective in producing the leaf curl deformation.

Figure 8 compares the intermediate result of the leaves wrinkle and curl deformation. From left to right, each column shows the intermediate deformation results of different leaves after iterating 0, 10, 20, 30, 40, 50, 60, 70, 80, 90, and 100 times. Obviously, as the iteration increases, the wrinkling and curling during the leaf withering process become more and more obvious, After 100 iterations, it can be seen that realistic leaves withering simulation results can be obtained with our approach.

5.3. Stress change of cells

Take the willow leaf model as an example, Figure 9 shows the change curve of each force magnitude on the cells at the feature points of the veins, vein points, boundaries, and mesophyll during the iterative deformation process, respectively.

As the iteration increases, It can be obviously seen that, firstly, the time-varying external forces defined at the feature points gradually increases, while under the



influence of the adjustment coefficient K_r , for cells in other regions, the increase in the external forces transmitted by the feature points of the leaf veins is not obvious.

Secondly, the elastic force increases in a wave-like way, which is caused by the tension or compression of the springs around the cells during the wrinkling process.

Thirdly, the cell swelling force increases steadily, which is because of the continued dehydration of cells, the internal pressure of the cell gradually decreases, and the difference between internal and external pressure increases constantly.

Finally, the damping force is directly proportional to the velocity of cells. The curve of damping force magnitude shows that the velocity of the leaves overall curling deformation is accelerating. But at the feature points of veins, the velocity decreases at the later stage of the iteration, which is because the leaf tips curl inward to make the cells decelerate in a certain direction at the feature points.

5.4. Cell swelling force and water loss rate

When dehydration occurs, the permeation effect of the cells reduces the water content, which makes the leaf curl and deforms under the cell swelling force. A significant factor to measure the water loss of cells is the water loss rate w_{loss} . As the iteration increases, Figure 10 compares the cell swelling force changing curves of a willow leafs

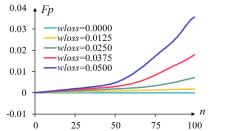


Fig. 10. Effect of water loss rate on cell swelling force

mesophyll cell with different water loss rates w_{loss} during the dehydration process.

Figure 11 compares the wrinkle deformation effect with different water loss rate w_{loss} after 100 iterations. It can be observed that when $w_{loss}=0$, i.e., when the cell is not dehydrated, there is no swelling force on the cell, and the wrinkle deformation doesn't exist. When $w_{loss} > 0$, which means as the number of iterations increases, the cell will continuously losing water. In this situation, with a larger water loss rate w_{loss} , the rate of the cell swelling force grows faster. As a result, cells have greater swelling force, and the effect of leaf wrinkle deformation is more pronounced.

5.5. The external force and coefficient

Figure 12 compares the effects of leaf curl deformation with different time-varying initial values of external forces

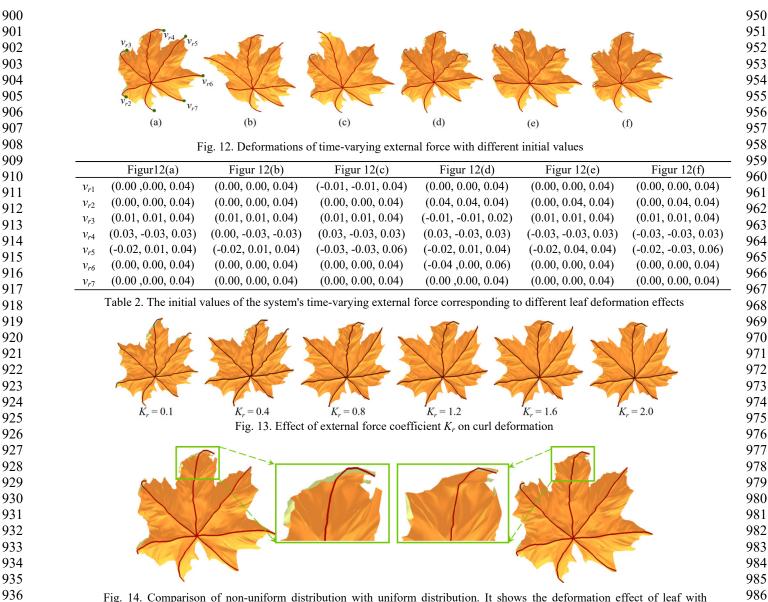


Fig. 14. Comparison of non-uniform distribution with uniform distribution. It shows the deformation effect of leaf with non-uniform mass distribution (a) and with uniform mass (d); (b) and (c) are the corresponding local amplification results.

(including size, direction and adjustment coefficient, etc.).Table 2 provides the details of corresponding values.Obviously, by adjusting the initial values of external forces, the algorithm we proposed in this article is well suited to generate a variety of curl deformation effects.

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943 With the same condition of external forces, different curl 944 deformation effects obtained by setting different external 945 force coefficient K_r can be compared in Figure 13, where K_r 946 is in the range [0, 2]. The comparison shows that when the 947 system time-varying external forces are constant, the extent 948 of curl deformation at the feature points of the leaf veins 949 was proportional to the external force K_r , while the range of curl deformation influenced by the system time-varying external force was inversely proportional to K_r .

5.6. Non-uniform mass distribution

In this article, we define a non-uniform mass distribution, 991 the cells that close to the veins have a larger mass, and the mass of ones that close to the edge and away from the veins are smaller. 992

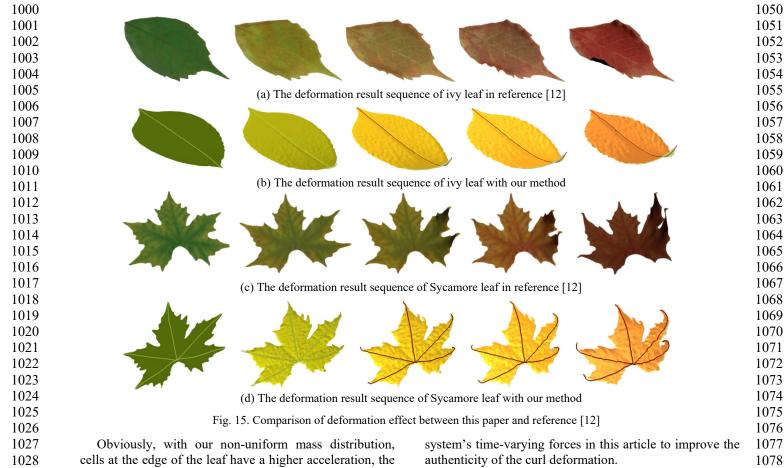
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Figure 14 compares the effect of the initial mass distribution of leaf cells on the wilting deformation. When the other influencing factors are the same, Figure 14(a) shows the wilting deformation result using our non-uniform mass distribution, and Figure 14(d) shows the result using uniform mass distribution. Figure 14(b) and Figure 14(c) are local amplifications of the corresponding effects.



corresponding degree of curl deformation is more pronounced.

5.7. Compare with related work

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Figure 15 shows the comparisons between the results of this article and reference [12]. From these comparisons, the following analyses can be obtained.

1036 Firstly, reference [12] constructs a two-layer model for 1037 the veins and a single-layer model for the mesophyll, 1038 which has an expensive cost in computation. In order to 1039 improve the efficiency, in this article, we adopt single-layer models for both of the veins and mesophylls 1040 1041 and use a non-uniform mass distribution to simulate the deformation. As a result, the veins we construct have 1042 higher quality and lower computational cost. 1043

Secondly, reference [12] only considers the elastic force 1044 and damping force in the simulation process. In contrast, in 1045 addition to the above two forces, the algorithm we proposed 1046 in this article also introduces the cell swelling force and 1047 defines a non-uniform mass distribution. It can therefore be 1048 assumed that the wrinkle deformation in this article is much 1049 more pronounced.

> Finally, it can be seen that the deformation simulations in reference [12] are less natural. In contrast, we introduce the

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6. Conclusion

1081 The main goal of the current study was to simulate the 1082 deformation of leaves wrinkle and curl due to the 1083 dehydration effect. To achieve this goal, we construct a 1084 cellular dynamics-based model of leaf wilt and curl 1085 deformation, which is based on the mass-spring model 1086 defined by cellular dynamics theory and combined with the swelling force describing cell permeation and the 1087 time-varying external forces. 1088

In this article, firstly, we propose a boundary expansion 1089 algorithm to locate the feature points of the veins and 1090 construct the primary veins based on discrete geodetic 1091 paths. Secondly, we adopt a mass-spring system of the cells 1092 according to the cell dynamics theory, and provide a 1093 non-uniform mass distribution to increase the acceleration 1094 of the movement of the mass cells in the boundary area. 1095 Thirdly, we propose a swelling force of the cells based on 1096 permeation theory and water content, and control this force 1097 by adjusting the water loss rate. Fourthly, we define the 1098 system's time-varying external forces on the veins, and 1099 adjust the initial and iterative parameters of the force to simulate the curl deformation of the leaf during the dehydration and wilting process. Finally, the equation of

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1101 motion is solved iteratively via the implicit midpoint 1102 method to get the simulation results.

Further research in this field would be of great help in 1103 improving the authenticity of the simulation. Firstly, in this 1104 article, we only consider the geometry simulation of leaves, 1105 and simply select color rendering without considering 1106 texture synthesis. Thus it would be interesting to introduce 1107 the texture mapping and color rendering process into the 1108 simulation process. Secondly, considerably more work will 1109 need to be done to modify the non-uniform mass 1110 distribution interpolation and improve the interpolation 1111 formula. Thirdly, It would be interesting to construct two or 1112 even multi-level veins to find the best compromise between 1113 simulation effect and computational efficiency. Finally, 1114 another interesting direction for future work is to take a 1115 large number of real leaf texture images with different 1116 degrees of wilting and combine them with the geometric 1117 deformation model for texture synthesis to further improve 1118 the visual effect of leaf wilting deformation from the color and texture perspective. 1119

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